

Topographic and orbital forcing of Titan's hydroclimate

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ABSTRACT

The cause of the hemispheric asymmetry of Titan's methane lakes and seas is the subject of ongoing debate. A leading hypothesis posits that seasonal insolation asymmetries caused by Saturn's eccentric orbit lead to differences in net precipitation over the two poles, perhaps mediated by asymmetric atmospheric transport of moisture. But topographic variations have also been proposed to contribute, albeit without considering the importance of surface hydrology. Here we present general circulation model simulations including a synchronously coupled surface and ground hydrology scheme, testing the separate and combined influences of topography and orbital forcing on Titan's hydroclimate. We find that, while topography leads to warmer polar regions relative to a flat surface which in turn enhances methane loss to the atmosphere, the overall effect on the global distribution of surface methane liquid is minor. In particular, topography does not force any notable asymmetry in the meridional circulation, nor does it affect the seasonality of the methane cycle, though it does increase the regional heterogeneity of average precipitation at mid-latitudes. We also find that Titan's atmospheric methane transport robustly responds to orbital forcing, in agreement with previous results, but this is insufficient to overcome the distribution of surface liquids dictated by surface hydrology. We conclude that Croll-Milankovitch cycles are plausible on Titan, but likely not the dominant driver of the current distribution of liquids; relatedly, our results suggest that the volume of the large seas and lakes has not varied substantially on millennial timescales.

1. Introduction

Titan supports an active hydrologic cycle, comprising observed features like clouds, rivers, lakes, and seas (Roe, 2012; Turtle et al., 2011a, 2018b; Hayes, 2016; Lopes et al., 2019), in which methane cycles through the atmosphere and crust on various timescales (Lunine and Lorenz, 2009; Roe, 2012; Mitchell and Lora, 2016; Hörst, 2017; Hayes et al., 2018). Titan's sporadic cloud activity evolves seasonally, and summer clouds have now been observed at the high latitudes of both hemispheres (Turtle et al., 2018b). These polar regions host surface lakes and seas of liquid methane (Hayes, 2016), some of whose shorelines have potentially retreated on sub-annual timescales (Turtle et al., 2011b; Hayes et al., 2011; MacKenzie et al., 2019a).

Separately, general circulation models (GCMs) of Titan have demonstrated that the climate system transports methane to the polar regions on multi-annual timescales, where liquid is cold-trapped in the surface reservoirs (Mitchell, 2008; Schneider et al., 2012; Lora et al., 2015; Newman et al., 2016). And, conversely, simulations with ample surface methane at the polar regions—beyond that available in the seas and lakes—best reproduce various aspects of the observed hydroclimate (Lora and Mitchell, 2015; Mitchell and Lora, 2016; Lora and Ádámkóvics, 2017; Faulk et al., 2017). (Geologic and remote sensing evidence similarly points to wet polar surfaces (e.g., Neish and

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41 Lorenz, 2014; Jennings et al., 2016.) Methane moisture, sourced at the poles, humidifies the lower latitudes via the
42 atmospheric circulation (Griffith et al., 2014; Mitchell and Lora, 2016; Lora and Ádámkovics, 2017; Battalio and Lora,
43 2021).

44 On longer timescales, orbital precession causes the seasonal insolation asymmetry due to Saturn's orbital
45 eccentricity to evolve on roughly 10^4 year timescales, and this has been proposed to contribute to or control the
46 asymmetric distribution of the surface liquids: Titan's large seas and the vast majority of small lakes occur in the north
47 (Aharonson et al., 2009; Hayes, 2016). Indeed, Lora et al. (2014) showed that a surface methane difference between
48 hemispheres in four paleoclimate GCM simulations corresponded to the hemispheric contrast in peak insolation
49 caused by changes in orbital parameters. Other studies have similarly found that the modern orbital configuration,
50 with perihelion occurring close to northern summer solstice, leads to a net accumulation of methane in the north
51 (Schneider et al., 2012; Lora et al., 2015; Lora and Mitchell, 2015; Newman et al., 2016). Lora and Mitchell (2015)
52 further showed that equatorward transport by atmospheric eddies is asymmetric, providing a mechanism to translate
53 the asymmetric seasonal forcing into a net transport by the atmosphere.

54 Surface observations of Titan, however, do not lead to an unequivocal interpretation in agreement with these
55 findings. Geomorphologic mapping of Titan's polar regions suggests that roughly equal areas in the two hemispheres
56 are made up of low-lying basins, which are filled in the north but empty in the south, suggestive of paleo-seas (Birch
57 et al., 2017, 2018). But putative evaporitic deposits, present extensively in the northern basins and shores of filled
58 lakes (as well as at lower latitudes), are conspicuously absent from the southern polar regions (MacKenzie et al.,
59 2014). Therefore, if methane cycles between poles on the timescale of precessional cycles, then evaporites left behind
60 in the south would need to have been buried or removed, implying that a complete emptying of southern lakes and seas
61 is not occurring at present and thus should have occurred relatively quickly in the recent past. A clear alternative is
62 that such liquids never existed in the past $\sim 10^4$ years, obviating the need for a net pole-to-pole transport mechanism,
63 and that basin formation was much more ancient.

64 Moreover, all of the above GCM studies used simplified approximations for surface hydrology and ignored the
65 influence of topography, disregarding, for example, the importance of runoff. Such simplifications could be important,
66 particularly given the substantial surface-atmosphere interactions and associated feedbacks. More recently, Tokano
67 (2019) used a GCM incorporating heterogeneous surface properties (topography, albedo, emissivity, and thermal
68 inertia) to suggest that an asymmetric meridional circulation, forced in particular by surface topography, hinders
69 precipitation in the southern hemisphere and is therefore the principal driver of any asymmetry. This would suggest
70 that Titan's lake distribution resists any influence of orbital forcing, and thus may not have changed significantly on
71 $<10^5$ year timescales.

72 Yet the study of Tokano (2019) also used a simplified hydrology scheme, neglecting lateral transport of surface
73 liquids; surface methane was exclusively affected by precipitation and evaporation. On the other hand, Faulk et al.
74 (2020) recently coupled a hydrology model and GCM, enabling for the first time a fully self-consistent simulation of
75 the methane cycle that demonstrated the importance of subsurface methane in enabling high-latitude surface moisture.
76 However, that model did not explicitly account for the direct influence of topography on the atmosphere, so evaluation
77 of such influence on the resulting meridional circulation was not possible.

78 The present study addresses that gap. Here, we present new simulations with the coupled GCM and hydrology
79 model (Faulk et al., 2020), wherein topography now interacts fully with the atmosphere. We evaluate the influence of
80 this interaction in the context of the hydroclimate including surface and subsurface hydrology. And, using our updated
81 model, we test the relative importance of orbitally-forced seasonal insolation asymmetries in the global transport of
82 methane, and the consequent distribution of lakes and seas.

83 **2. Methods**

84 The model we use in this study is a configuration of the Titan Atmospheric Model (TAM; Lora et al., 2015).
85 TAM is a three-dimensional GCM based on the GFDL Flexible Modeling System, which, in addition to its spectral
86 dynamical core, uses component modules to parameterize unresolved physical processes including moist convection,
87 large-scale condensation, surface heat conduction, surface–atmosphere fluxes, turbulent diffusion, and full radiative
88 transfer (Lora et al., 2015; Lora and Ádámkóvics, 2017; Lora et al., 2019). We run the model at T21 resolution (64
89 longitude by 32 latitude grid points), with 24 unevenly spaced vertical levels, and a time-step of 600 s. In all cases, we
90 use a constant surface thermal inertia of $750 \text{ J m}^{-2} \text{ s}^{0.5} \text{ K}^{-1}$ (MacKenzie et al., 2019b). In addition, we employ several
91 recent updates including a tuned convection parameterization and other incremental improvements (Battalio et al.,
92 2021) and, crucially, a model of surface and subsurface methane hydrology that enables the self-consistent simulation
93 of the methane cycle (Faulk et al., 2020).

94 The hydrology model, which calculates infiltration, runoff, and a subsurface methane table with lateral flow, uses
95 topography as an input (Faulk et al., 2020). Here, we additionally incorporate the direct effects of topography on
96 Titan’s atmosphere by including it as a lower boundary condition in selected simulations. To accomplish this, we use
97 the geoid-corrected map of Titan’s topography estimated from Cassini RADAR data (Corlies et al., 2017), slightly
98 modified to enforce a flat equipotential surface over Titan’s maria. In addition, we do not regularize the topography.
99 For consistency, we use the same topography as a boundary condition for both the atmosphere and for the hydrology
100 model, and as a result the latter is somewhat different than that in Faulk et al. (2020); the large-scale features are
101 nevertheless equivalent.

102 For this work, we run a number of simulations to investigate the individual and cumulative effects of the topography
103 and of orbital forcing variations on Titan’s hydroclimate. First, we run four simulations that use 1) either a flat surface (as
104 in previous studies) or the new topography, and 2) either Saturn’s modern orbital parameters or a case with the longitude
105 of perihelion shifted by 180° with respect to modern. (This latter forcing causes the shorter, more intense summer to
106 occur in the northern hemisphere, while other orbital parameters are held constant; this is slightly different than the
107 approach used by Lora et al. (2014), but roughly corresponds to their 28-kyr-before-present case.) The simulations
108 are otherwise identical. We note, in particular, that in all cases the hydrology model uses the input topography as a
109 prerequisite for calculating runoff and seepage; in other words, only the lower atmospheric boundary condition is flat
110 in the “flat surface” simulations, as in Faulk et al. (2020).

111 In all four of these cases, we use a value for the hydraulic conductivity of the surface $k = 5 \times 10^{-5} \text{ m s}^{-1}$ within
112 the hydrology model based on previous results with a flat surface (Faulk et al., 2020). We find, however, that while
113 our flat simulations behave as expected, the simulations using topography begin to lose their surface liquid methane
114 after about four decades, and thus do not trend toward a statistical equilibrium that corresponds to Titan’s observed
115 surface. As a result, we run two additional simulations (one each for the two longitude of perihelion configurations)
116 with topography and with a higher hydraulic conductivity $k = 10^{-4} \text{ m s}^{-1}$. These are restarted from the end of the
117 40th Titan year of the corresponding lower- k simulations with topography.

118 In all cases, simulations are run out to the 100th year (henceforth, references to “years” and “decades” correspond to
119 Titan years), and all comparisons between simulations that follow are done using corresponding years (unless otherwise
120 noted, figures show results from the final decade of simulations). Simulations are initialized at year zero using no
121 surface liquid and a prescribed amount of subsurface methane, as in Faulk et al. (2020), as well as with a previously
122 spun-up atmospheric state to minimize computational expense. All simulations reach statistical steady state for relevant
123 atmospheric variables, like temperature, winds, and humidity, within 20 years. While there is still some interannual
124 variability in, for example, the total surface liquid by the 100th year, we find that our results are robust when using
125 multiple years (and this allows an estimate of uncertainty due to variability). The exception, of course, is the pair of
126 simulations (with topography and $k = 5 \times 10^{-5} \text{ m s}^{-1}$) mentioned above, which consistently lose surface methane in
127 the final several decades.

128 **3. Results**

129 **3.1. Impact of topography**

130 We begin by investigating the impact of topography as a boundary condition for the atmosphere in the simulations
131 using Titan’s modern effective orbital configuration (in other words, Saturn’s orbit around the Sun). The principal
132 questions motivating this are whether the topography introduces any substantial structure into hydroclimate patterns,

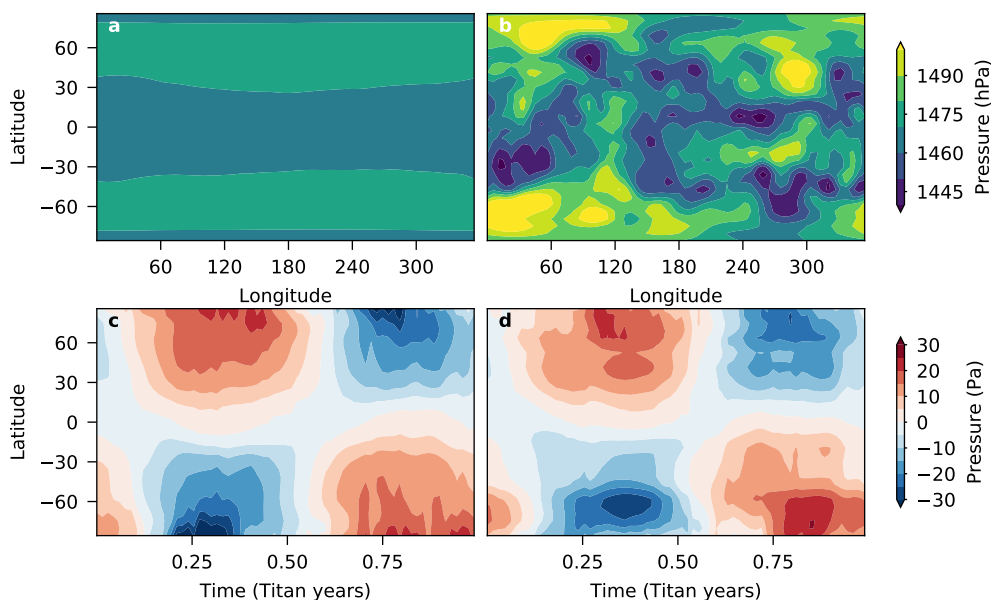


Figure 1: Surface pressures. Top: time-averaged surface pressures for simulations (a) without and (b) with topography. Bottom: climatological seasonal evolution of zonally averaged anomalous surface pressure for simulations (c) without and (d) with topography. The beginning of the year corresponds to northern autumnal equinox.

133 and whether or not it instigates a significantly asymmetric Hadley circulation as it does on Mars (e.g., Richardson and
 134 Wilson, 2002; Zalucha et al., 2010).

135 The annually-averaged surface pressure, as well as its climatological (i.e., the decadal average) seasonal anomalies,
 136 from simulations with and without topography, are shown in Fig. 1. Immediately apparent is the fact that topography
 137 leads to large variations of the average surface pressure—on the order of tens of hPa—on regional scales, corresponding
 138 to the distribution of valleys and ridges of the surface. Indeed, because the low-lying basins cluster at higher latitudes,
 139 the surface pressure tends to be higher near the poles and lower over the low-latitude highlands, with the exception
 140 of the region of Xanadu (roughly 20°S, 240–300°E), which is thought to be relatively low-lying, albeit mountainous
 141 (Lorenz et al., 2013; Corlies et al., 2017). All of this is in stark contrast to the case without topography, wherein the
 142 average surface pressure variations are very weak.

143 On the other hand, seasonal surface pressure anomalies are small for both simulations (on the order of tens of
 144 Pa between summer and winter at the poles), consistent with Titan’s weak Coriolis force and temperature gradients
 145 (Mitchell and Lora, 2016). Importantly, these anomalies are also effectively indistinguishable between simulations,
 146 indicating that topography has little meaningful impact on the movement of mass in Titan’s atmosphere.

147 Consistent with this, the climatological seasonal anomalies of surface temperature (Fig. 2a, d) and column-
 148 integrated methane (Fig. 2b, e) are also largely interchangeable between simulations, showing no appreciable difference

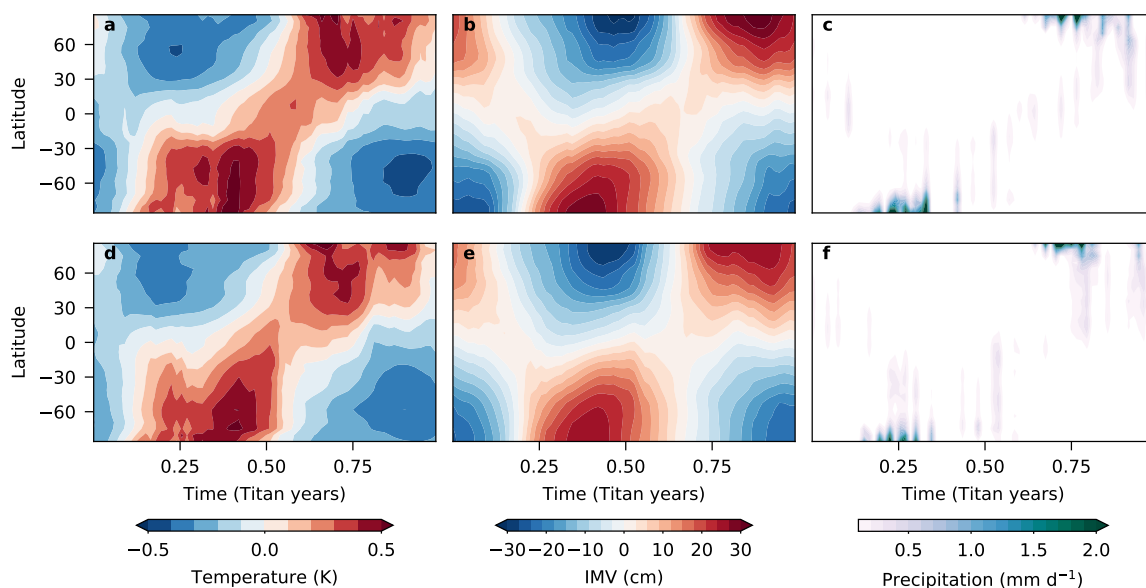


Figure 2: Seasonal variations. Climatological seasonal evolution of zonally averaged (a, d) anomalous surface temperatures, (b, e) anomalous column-integrated atmospheric methane vapor (IMV), and (c, f) precipitation, for simulations (a–c) without and (d–f) with topography (in f, the case with $k = 10^{-4} \text{ m s}^{-1}$ is shown). The beginning of the year corresponds to northern autumnal equinox.

149 either in magnitude or timing. These diagnostics indicate that, while topography affects the time-averaged distribution
 150 of atmospheric mass (as expected, since there is less atmosphere over a tall mountain, for instance) and correspondingly
 151 the column methane vapor at any one location, it does little to modify Titan’s seasonal circulation patterns and humidity
 152 variations relative to the well-studied case without topography.

153 Comparison of the seasonality of zonally-averaged precipitation between analogous simulations with topography
 154 (that is, simulations in which the methane cycle is equally active; see Section 2) similarly shows almost no differences
 155 (Fig. 2c, f): precipitation peaks over the summer poles and occurs sporadically at other latitudes, with a slight but
 156 consistent enhancement at northern relative to southern latitudes during their respective summers. The single notable
 157 difference in precipitation distributions is the modest increase in southern midlatitude activity in early summer in the
 158 simulation with topography, which could be a better match to early observations of Titan’s clouds (e.g., Roe et al.,
 159 2005) and will be further discussed below.

160 Figure 3 shows the climatological seasonal meridional mass stream function in Titan’s lower atmosphere for the
 161 two simulations. As in previous studies, the model simulates a seasonal Hadley circulation with relatively strong,
 162 cross-equatorial winter cells and weak, shallow summer cells, accompanied by relatively steady and shallow thermally
 163 indirect cells at higher latitudes, which are associated with eddy activity (Lora and Mitchell, 2015). While quantitative

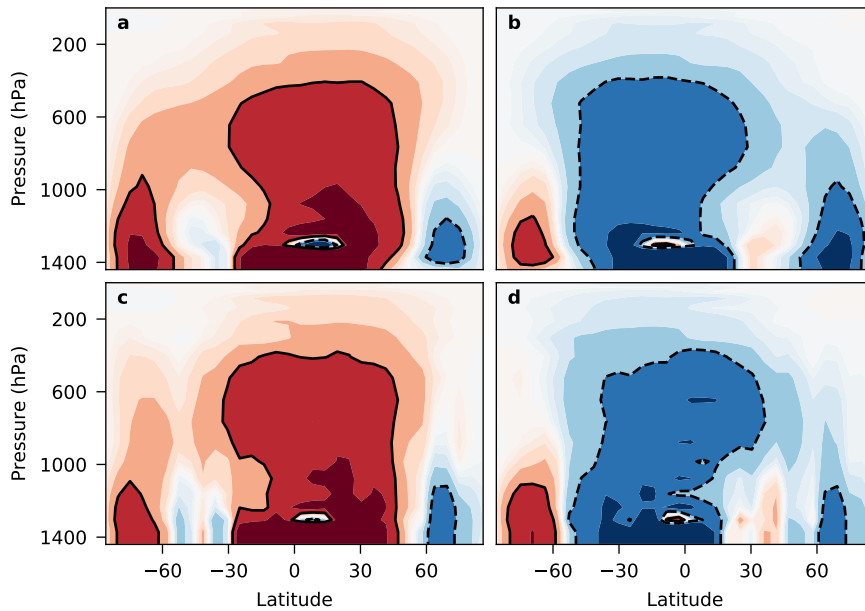


Figure 3: Meridional mass stream function. Time-averaged meridional mass stream functions over northern fall and winter (a, c) and northern spring and summer (b, d) for simulations (a, b) without and (c, d) with topography. Positive values indicate clockwise motion. Contours increase by factors of 2; the $\pm 0.5 \times 10^9 \text{ kg s}^{-1}$ contours are shown in solid (positive) and dashed (negative) black contours.

164 differences exist, there is little indication of significant qualitative differences in the meridional circulation from the
 165 simulation with topography relative to the one without. In particular, we find no evidence of a hemispheric asymmetry
 166 induced by the topography. We therefore discard this as a plausible mechanism to drive the asymmetry of surface
 167 liquids on Titan.

168 Importantly, all of the above results do not mean that topography has *no* effect on Titan's hydroclimate. Consistent
 169 with the distribution of time-averaged surface pressures, the distribution of time-averaged surface temperatures is
 170 affected by the topography (Fig. 4): lower elevation surfaces are warmer than other regions as a simple consequence of
 171 the lapse rate. But because this principally affects higher latitudes, the ultimate result is that the equator-to-pole surface
 172 temperature difference is approximately 0.6 K smaller on average in the case with topography. This, in turn, means
 173 that the poles are weaker cold traps for surface liquid, which initially promotes a more vigorous methane cycle but, as
 174 described in the previous section, eventually leads to a loss of surface methane in the case with $k = 5 \times 10^{-5} \text{ m s}^{-1}$.
 175 Less surface and near-surface liquid leads to lower evaporation rates, which in turn causes higher surface temperatures
 176 (Mitchell, 2008; Lora et al., 2015; Lora and Ádámkóvics, 2017; MacKenzie et al., 2019b); the result, after 100 years,
 177 is shown in Fig. 4b.

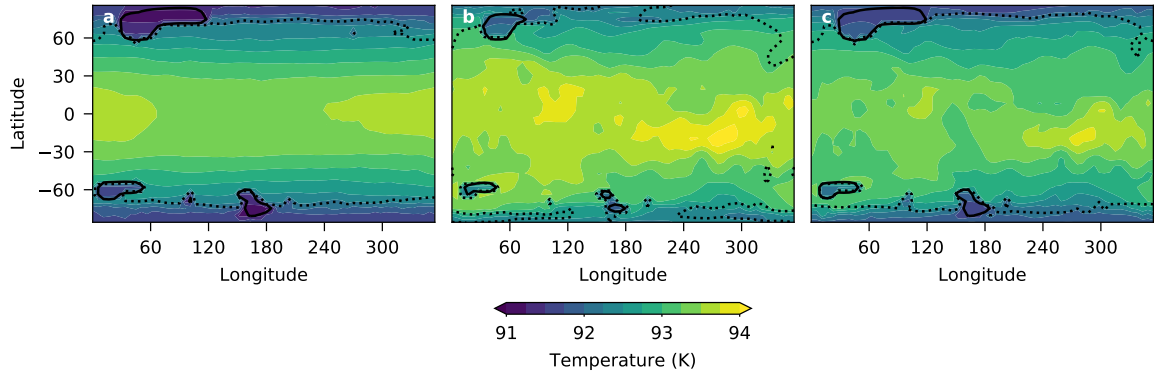


Figure 4: Surface temperatures. Time-averaged surface temperatures for simulations (a) without topography, (b) with topography, and (c) with topography and increased hydraulic conductivity. Locations with surface liquid methane exceeding 25 m are outlined in black solid contours; locations where average evaporation exceeds 0.1 mm d^{-1} (“day” refers to Earth days) are outlined in black dotted contours.

178 Figure 4c shows the equivalent time-averaged surface temperatures for the simulation with topography and
 179 $k = 10^{-4} \text{ m s}^{-1}$. In this case, the higher hydraulic conductivity allows seepage of subsurface methane to offset
 180 atmospheric divergence of methane moisture from the high latitude surface sources, balancing the evaporation rates
 181 and allowing for a stable hydroclimate in the last several decades of simulation (see Fig. 2f). Under this configuration,
 182 the global average surface temperature is similar to the case without topography (Fig. 4), though still equator-to-pole
 183 contrasts are weaker, and the polar cold traps are therefore also weaker (Fig. 4a, c).

184 The corresponding distributions of surface and near-surface liquids for the final decade of the simulations are also
 185 illustrated in Fig. 4 (near-surface liquid is indicated by regions with substantial average evaporation rates). Interestingly,
 186 the total volume of surface liquids at the end of our simulations is on the order of 10^5 km^3 , in excellent agreement with
 187 the total volume of Titan’s liquids estimated from observations (Lorenz et al., 2008; Mastrogiuseppe et al., 2014).

188 In all cases, surface liquids pool in areas of topographic lows corresponding to Titan’s Kraken and Ligeia maria
 189 and the catchment basin of Ontario Lacus, as well as in the lowest-elevation basin in the southern hemisphere (around
 190 60°S , 30°E), a location where liquids have not actually been observed. All of this is primarily due to the routing of
 191 runoff, in combination with the latitudinal balance of precipitation and evaporation (Faulk et al., 2020). Outside of
 192 these surface liquids, ample evaporation (from near-surface ground methane) occurs at high latitudes, corresponding
 193 to the subsurface methane reservoirs and moist regolith inferred in previous works (Neish and Lorenz, 2014; Lora and
 194 Mitchell, 2015; Mitchell and Lora, 2016; Jennings et al., 2016; Lora and Ádámkóvics, 2017; Turtle et al., 2018b; Faulk
 195 et al., 2020).

196 In the simulation with topography and $k = 5 \times 10^{-5} \text{ m s}^{-1}$, surface liquids and high evaporation rates cover
 197 substantially smaller areas than in the flat case by the end of the simulations (Fig. 4b), consistent with a foundering

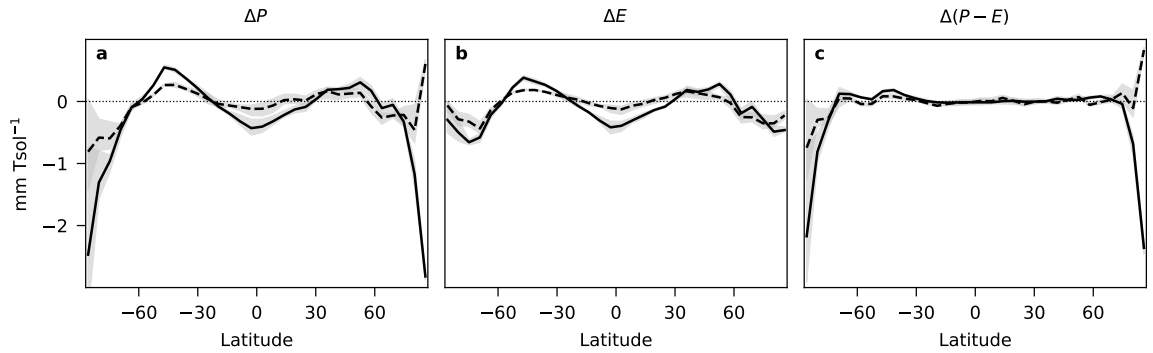


Figure 5: Responses of surface moisture fluxes to topography. Time- and zonal-mean differences in (a) precipitation, (b) evaporation, and (c) precipitation minus evaporation between simulations with and without topography (solid curves), as well as between simulations with topography plus increased hydraulic conductivity and without topography (dashed curves). In all cases, the gray shading shows the range of decadal averages over the final three decades of the simulations.

198 methane cycle. On the other hand, in the simulation with topography and $k = 10^{-4} \text{ m s}^{-1}$, Kraken/Ligeia Mare and
 199 Ontario Lacus, as well as regions with near-surface ground methane, are similar in area as in the flat case, while the
 200 other (unobserved) southern reservoir is smaller, therefore better approximating observations (Fig. 4c).

201 These points are further demonstrated by a quantification of precipitation and evaporation differences between the
 202 simulations. Fig. 5 shows the impacts of topography, as well as an estimate of uncertainty due to variability, from the
 203 final three decades of the simulations. Under the same hydraulic conductivity, topography leads to modest increases in
 204 average precipitation in the mid-latitudes, as well as a modest decline near the equator; the largest impact, however, is in
 205 substantial declines at higher latitudes, which approach 2 mm Tsol^{-1} (“Tsol” refers to a Titan day) on average near the
 206 poles (Fig. 5a). With higher hydraulic conductivity, these differences are muted and, in particular, polar precipitation
 207 differences are smaller and of opposite signs in the north and south.

208 The evaporation response closely follows that of precipitation (Fig. 5b), since precipitation enables subsequent
 209 evaporation from lower latitudes, which are otherwise relatively dry. The exception is, expectedly, at higher latitudes,
 210 where ground methane is sufficiently close to the surface to provide a source of liquid (as illustrated in Fig. 4).
 211 Evaporation differences between the cases with and without topography suggest less polar evaporation in the former,
 212 as a result of the evaporation line being farther poleward; this is especially true for the case with lower hydraulic
 213 conductivity, as discussed above.

214 In the net, differences in surface–atmosphere fluxes of moisture largely cancel out at most latitudes, with the
 215 exception of the polar regions (Fig. 5c). There, the simulation with topography and lower hydraulic conductivity clearly
 216 has less net precipitation ($P - E$) than the flat case, which leads to the loss of surface methane described above. On
 217 the other hand, the simulation with topography and higher hydraulic conductivity has lower net precipitation in the
 218 south, but higher net precipitation in the north, than the flat case. In other words, this simulation has a slightly enhanced

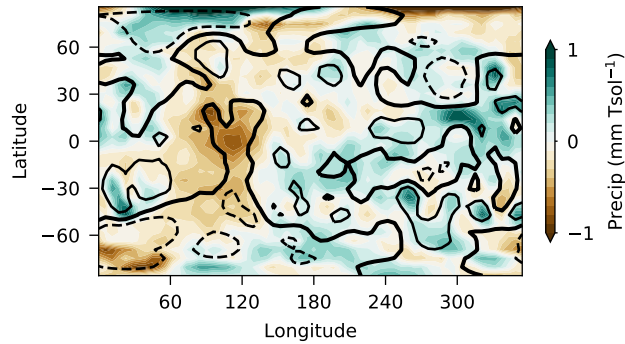


Figure 6: Zonally anomalous precipitation differences. Time-mean differences in the zonally anomalous (with the zonal average removed) precipitation between simulations with topography plus increased hydraulic conductivity and without topography, over the final three decades of the simulations. Topography is indicated by black contours, with solid and dashed contours indicating values 300 m above and below the average, which is shown in bold.

hemispheric asymmetry in net precipitation due to topography, although, as seen in Fig. 4, this is only barely reflected in the surface liquids. We note that, on average, the southern polar surface has slightly lower elevations than does the northern, and at the same time the deepest basin occurs more equatorward than do the deep basins of the north. In combination, these properties result in a modestly weaker cold trap and therefore a strengthened north–south net precipitation asymmetry, but only in the case where hydraulic conductivity is high enough that polar surface liquids are stable over long periods. These impacts do not depend on asymmetries of the meridional circulation.

A distinct impact of the topography is that it introduces inhomogeneities on regional scales. Southern mid-latitude precipitation is enhanced relative to the flat case, as can be seen in Fig. 5a, and this is accompanied by a slight depression of precipitation at the equator. While the absolute magnitudes are modest, these changes correspond to a threefold enhancement at southern mid-latitudes and a nearly 100% depression at the lowest elevations near the equator. This reflects both the low totals in these regions, but also the important influence of topography on the occasional rainstorms that affect them. Indeed, at mid- and lower latitudes, regions of relatively low (high) precipitation occur in regions of relatively low (high) elevation (Fig. 6); this is particularly obvious at the equator and 120°E, as well as in the highlands surrounding Xanadu. Especially over the topographic rises around 40°S and between 180 and 330°E, precipitation is enhanced; these results constitute indications of orographic precipitation on Titan (see Barth, 2010), which merit further study. A possible geographic control of southern midlatitude clouds was suggested based on cloud observations early in Titan’s southern summer (Roe et al., 2005), although this was later discounted. Our results here suggest that topography does in fact locally enhance precipitation (and, by extension, precipitating clouds, though we do not explicitly model cloudiness), although precipitation at the same latitudes also occurs at other longitudes. In short, precipitation is enhanced (depressed) over high (low) elevation and increases at mid-latitudes (in both hemispheres) as a result of topography, but none of this has obvious consequence on the distribution of surface liquids.

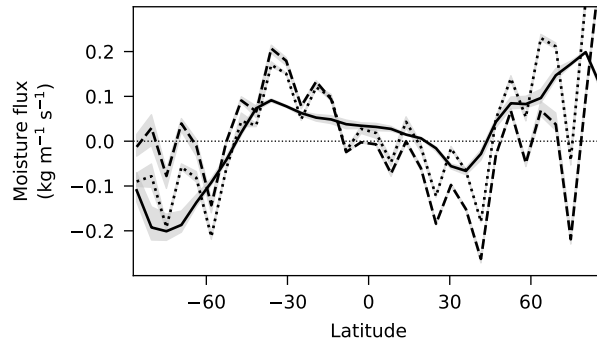


Figure 7: Meridional moisture fluxes. Time- and zonal-mean atmospheric fluxes of moisture from simulations without topography (solid curve), with topography (dashed curve), and with topography and increased hydraulic conductivity (dotted curve). In all cases, the gray shading shows the range of decadal averages over the final three decades of the simulations. Positive values indicate northward transport.

240 3.2. Impact of orbital forcing

241 Having established that the latitudinal distribution of surface liquids in our simulations with Titan’s modern orbital
 242 configuration is largely determined by surface hydrology and only indirectly affected by the influence of topography
 243 on the atmosphere, we next turn to assessing the importance of orbital forcing on Titan’s hydroclimate in the current
 244 version of TAM. The primary aim is to understand whether seasonal asymmetries due to the eccentricity of Saturn’s
 245 orbit produce a first- or second- order effect in our atmosphere–surface coupled model.

246 As a starting point, the net atmospheric moisture fluxes from the three simulations with modern orbital configura-
 247 tion are shown in Fig. 7. In the case without topography, this consists of poleward moisture transport at high latitudes
 248 and equatorward transport at mid-latitudes, with a hemispheric asymmetry that results in cross-equatorial transport
 249 and convergence in the northern hemisphere; these results are quite similar to the idealized case described in Lora
 250 and Mitchell (2015), including in their magnitude. It is worth emphasizing that these represent the net transport that
 251 comprises the residual of much larger, mostly opposing cross-equatorial seasonal fluxes (Lora and Mitchell, 2015).

252 Topography introduces considerable noise into the net moisture flux curves (Fig. 7), which is due to the zonal
 253 averaging that also results in residuals from opposing fluxes. Nevertheless, the overall global signal is still apparent; the
 254 case with $k = 10^{-4} \text{ m s}^{-1}$ (dotted line) largely resembles the flat simulation, whereas the case with $k = 5 \times 10^{-5} \text{ m s}^{-1}$
 255 (dashed line) produces generally more equatorward transport everywhere. In both cases, hemispheric asymmetries are
 256 difficult to see, though poleward transport is somewhat stronger at northern than southern high latitudes.

257 The total integrated northward transport of moisture in each simulation is shown in Fig. 8. With this metric,
 258 the impact of orbital forcing is more apparent, and the net atmospheric transport in the simulations coincides with
 259 expectations: northward in the modern, southward with a reversed perihelion, with approximately the same magnitude.
 260 However, this is only true when surface liquids are sufficiently stable (and available to evaporate): in the cases with

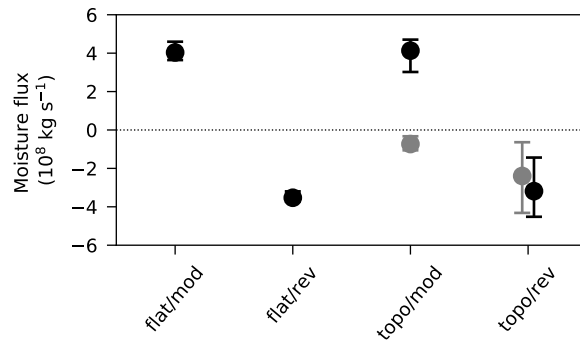


Figure 8: Net moisture transport. Net atmospheric moisture transport in each simulation: “Flat/topo” refers to simulations without or with topography, and “mod/rev” refers to the perihelion case, either modern or reversed. Gray symbols denote the results for the simulations with topography and lower hydraulic conductivity in which the surface liquids steadily decrease. Both the average and range of decadal averages over the final three decades of the simulations are shown. Positive values indicate northward transport.

261 topography and lower hydraulic conductivity (and disappearing surface liquids), the modern net transport is about zero.
 262 This latter point is an indication of the dependence of the orbital forcing mechanism (e.g., Aharonson et al., 2009; Lora
 263 et al., 2014; Lora and Mitchell, 2015) on the actual availability of surface liquids. In other words, without sufficient
 264 methane to evaporate into the atmosphere, asymmetric atmospheric transport cannot occur.

265 The total atmospheric moisture transport is significant, and represents, per Titan year, a northward (or southward
 266 under reversed perihelion) mass flux of methane approximately one order of magnitude larger than the mass present in
 267 the surface liquids. Of course, this does not correspond to the actual pole-to-pole mass transport, and instead mainly
 268 reflects the fact that most of the observable methane on Titan is in the atmosphere. Nevertheless, the atmospheric
 269 moisture flux clearly responds to the seasonal asymmetry induced by the orbital configuration.

270 On the other hand, the latitudinal distributions of differences in atmospheric moisture transport produced by
 271 simulations with a reversed longitude of perihelion provide a less consistent picture (Fig. 9). In all cases, these
 272 differences are negative at lower latitudes, indicating more southward net transport in the reversed perihelion case,
 273 in agreement with the reversed moisture flux asymmetry with reversed seasonal asymmetries. However, the transport
 274 differences at the poles are more complicated, again indicating their dependence on the details of the actual surface
 275 liquid distribution in each case. Notably, there is either a marginal (Fig. 9a,b) or robust (Fig. 9c) *northward* transport
 276 increase at the south pole, suggesting slightly increased *divergence* of methane from the south polar regions under
 277 reversed perihelion. In other words, our simulations seem to imply that atmospheric moisture transport out of, rather
 278 than into, the southern polar regions would increase if the seasonal asymmetry were reversed. Again, this could reflect
 279 the importance of the latitudinal distribution of polar basins, which is different in the two hemispheres.

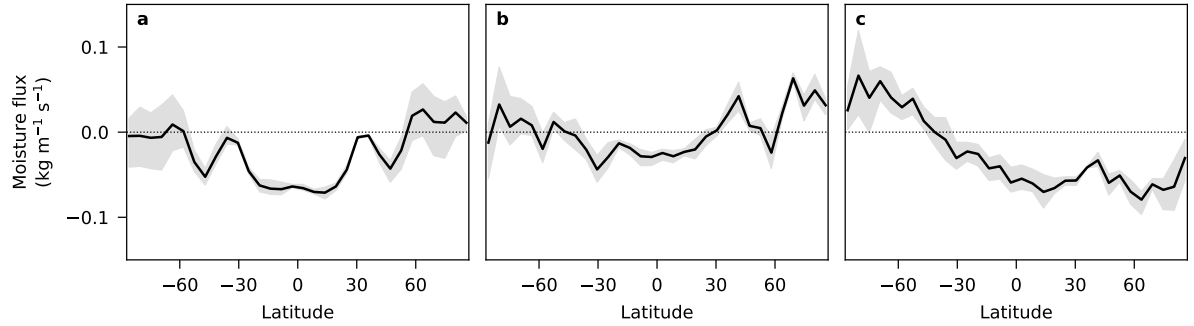


Figure 9: Moisture flux differences. Time- and zonal-mean differences in meridional atmospheric moisture flux between the reversed perihelion and modern simulations (a) without topography, (b) with topography, and (c) with topography and increased hydraulic conductivity. In all cases, the gray shading shows the range of decadal averages over the final three decades of the simulations. Positive values indicate northward transport.

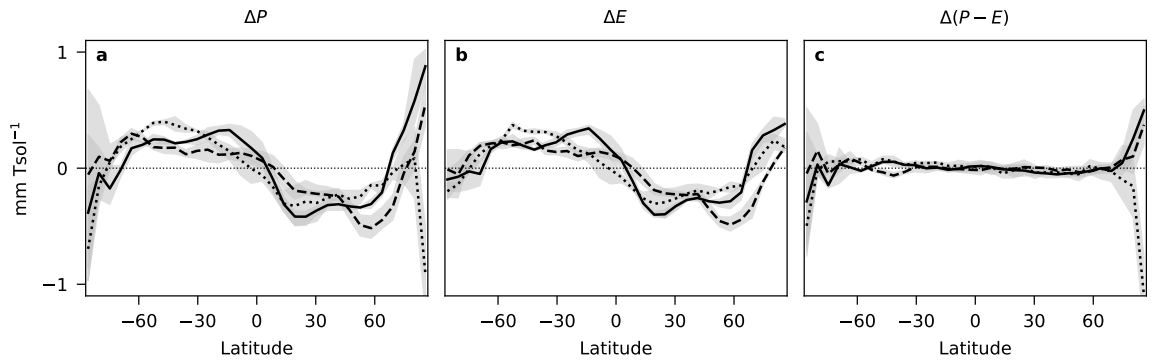


Figure 10: Responses of surface fluxes to orbital forcing. Time- and zonal-mean differences in (a) precipitation, (b) evaporation, and (c) precipitation minus evaporation between reversed perihelion and modern simulations. Results are shown for simulations without topography (solid curves), simulations with topography (dashed curves), and simulations with topography and increased hydraulic conductivity (dotted curves). In all cases, the gray shading shows the range of decadal averages over the final three decades of the simulations.

280 Finally, the precipitation and evaporation responses to reversed perihelion are shown in Fig. 10. In all cases, the
 281 global-scale response is the same: reversed perihelion leads to increased (decreased) precipitation in the southern
 282 (northern) hemisphere (Fig. 10a), with subtle latitudinal differences that imply a weak dependence on the topography as
 283 well as other factors, like the surface liquids, since the simulations with topography but different hydraulic conductivity
 284 produce slightly different results. The orbital forcing signal is less clear at the poles, perhaps indicating that stronger
 285 cold trapping leads to lower precipitation, but strongly modulated by surface heterogeneity. Notably, precipitation rates
 286 over the southern pole are *lower* in the reversed perihelion simulations in all cases, albeit with considerable variability.
 287 In addition, the precipitation responses are not very strong: in terms of methane mass, the differences in precipitation
 288 are largest at lower latitudes, peaking at magnitudes of 10^{14} kg yr⁻¹ around 15–45° in both hemispheres; this is orders
 289 of magnitude less than total annual precipitation.

290 As with the responses to topography, the evaporation responses to orbital forcing look very similar to those of the
291 precipitation, except at the poles where evaporation from subsurface methane is important (Fig. 10b). The result is
292 that the net precipitation ($P - E$) response to orbital changes is almost zero everywhere, with noisy responses at the
293 poles (Fig. 10c). Only the case with topography and higher hydraulic conductivity suggests less net precipitation at
294 the northern pole, and then without an accompanying increase in the south. Therefore, while net atmospheric moisture
295 transport robustly reverses with a reversed perihelion (Fig. 8), orbital forcing does *not* produce a clear signal in net
296 precipitation.

297 In fact, partly as a consequence of the above, surface liquids in our simulations do not show a consistent response
298 to the orbital forcing. This further indicates the importance of the interaction and feedbacks between mechanisms that
299 lead to the distribution of liquids, and their responses to seasonal asymmetries. There is not a simple precipitation or
300 evaporation response in our coupled model that can neatly explain the observed asymmetry of Titan's surface liquids.
301 We conclude that such an asymmetry is likely not a simple expression of orbital forcing, despite the clear atmospheric
302 response to changes in seasonal asymmetries.

303 4. Discussion and Conclusions

304 With our coupled, self-consistent model of Titan's climate and methane cycle, we have investigated the influence
305 of topography and orbital forcing on the hydroclimate with a series of simulations. Our results show that many
306 mechanisms can contribute to the distribution of surface liquids, and thus to the observed asymmetry of these liquids,
307 in our model. In particular,

- 308 • runoff and subsurface hydrology, and the corresponding distribution of basins (a result of topography), are
309 important drivers of the distribution of surface liquids at high latitudes;
- 310 • the impact of topography on surface temperatures, as a consequence of the lapse rate, results in the lower-
311 elevation polar regions being warmer, which promotes evaporation and thus, all else being equal, relatively
312 depletes liquids and can eventually throttle the methane cycle;
- 313 • relatedly, the rate of replenishment of evaporated liquids from ground methane helps determine whether certain
314 regions dry out or not, but depends, among other things, on the (unknown) hydraulic conductivity;
- 315 • topography affects the distribution of precipitation, particularly enhancing mid-latitude precipitation, but only
316 weakly impacts north–south hemispheric contrasts;
- 317 • orbital forcing produces a robust response in net atmospheric meridional moisture transport, in agreement with
318 previous results;

319 • but this transport is also sensitive to other factors and does not unequivocally lead to a strong response of surface
320 liquid distribution to orbital forcing.

321 Obviously, we do not yet know the details of the hydraulic conductivity of Titan's regolith or the extent and
322 connectivity of the putative methane table, and our results suggest (as in Faulk et al., 2020; Horvath et al., 2016)
323 that these details are important. For example, if the southern pole has lower hydraulic conductivity than does the
324 northern pole, that alone could potentially explain the observed asymmetry, despite the fact that many mechanisms
325 exist that could contribute, as described above. Needless to say, it is possible that other surface heterogeneities or
326 other complications not considered here, as well as model biases, could also play a role. Our simulations qualitatively
327 reproduce the modern methane distribution, but this could be associated with an over-constrained model, which would
328 limit its ability to accurately capture other climate equilibria. What is clear is that any signature of Titan's paleoclimate
329 recorded on the surface may be more difficult to interpret than previously thought, given all of these interacting
330 uncertainties.

331 Our results demonstrate that Titan's topography plays two major roles in the climate system—through its direct
332 interaction with the atmosphere, and by determining runoff directions and the distribution of basins—but neither
333 depends on significantly asymmetric meridional circulation cells, as is the case, for example, on Mars. At the same
334 time, our results indicate that no glaring inconsistencies with our current understanding of the Titan climate system
335 are introduced by using topography as a boundary condition, lending some further validity to the existing estimates
336 (Corlies et al., 2017). Naturally, improved knowledge of Titan's real topography is essential for future progress, as it
337 would eliminate ambiguities regarding the presence or absence of individual geographical features and enable more
338 targeted regional studies.

339 Separately, the question of whether or not the northern seas are growing in the present (or shrank in the past) as a
340 consequence of the orbitally-paced global climate remains unresolved. Despite the fact that the atmospheric circulation
341 does, in the net, transport methane northward in the current epoch (as long as there is a sufficient source in the south),
342 this net transport does not cleanly result in surface liquid volume increases in the north relative to the case with reversed
343 perihelion in our model. In all cases examined here, more surface liquid accumulates in the northern hemisphere than
344 the southern by the end of our simulations. This is in agreement with the results of Tokano (2019), albeit for starkly
345 different reasons.

346 Similarly, only in the simulations with topography and increased hydraulic conductivity is the north–south
347 asymmetry of surface liquids smaller for the reversed perihelion case than for the modern, and even then this is a
348 consequence of more relative loss of the northern seas. In none of our simulations does a reversed perihelion lead
349 to larger southern lakes. If our simulations are representative of Titan, then the southern basins are a relic of a past
350 climate beyond the recent Croll-Milankovitch cycles (Birch et al., 2018), suggesting that Ontario Lacus has not changed

351 in size significantly in the past tens of millennia (Tokano, 2021) and explaining the lack of evaporites at the south pole
352 (MacKenzie et al., 2014). It does not, however, preclude modest evolution of lake and sea shorelines, which could only
353 be captured by higher resolution or regional models. Thus, whether or not Titan's surface holds intelligible records of
354 millennial climate change remains an open question.

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